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Pentagonal Planar AX₅ Species: Synthesis and Characterization of the Iodine (III) Pentafluoride Dianion, IF₅^{2-v}

Karl O. Christe, *,‡,§ William W. Wilson, ‡ Greg W. Drake, # David A. Dixon, I Jerry A. Boatz, and Robert Z. Gnann§

Hughes STX, Propulsion Directorate, Air Force Research Laboratory, Edwards Air Force Base, California 93524, Loker Hydrocarbon Research Institute, University of Southern California, Los Angeles, California 90089, and Pacific Northwest National Laboratory, Richland, Washington 99352.

Abstract: The new $\mathrm{IF_5}^2$ dianion, which is only the second known example of a pentagonal planar $\mathrm{AX_5}$ species, was prepared as its $\mathrm{N(CH_3)_4}^+$ salt from $\mathrm{N(CH_3)_4}\mathrm{IF_4}$ and $\mathrm{N(CH_3)_4}\mathrm{F}$ in $\mathrm{CH_3CN}$ solution. Its structure was established by infrared and Raman spectroscopy, *ab initio* calculations and a comparison to isoelectronic $\mathrm{XeF_5}$. Furthermore, vibrational spectroscopy and x-ray powder diffraction data show that the previously reported composition "Cs₃IF₆" is actually a mixture of $\mathrm{Cs_2IF_5}$ and CsF. Ab inito calculations also show that the most probable geometries for $\mathrm{IF_6}^3$ are vibrationally unstable and undergo spontaneous F ion loss with formation of either pentagonal planar $\mathrm{IF_5}^2$ or square planar IF₄. The synthesis and some properties of the new $\mathrm{N(CH_3)_4IF_4}$ salt and a revised normal coordinate analysis of $\mathrm{XeF_5}$ are also presented.

Introduction

In 1991, the synthesis and characterization of the XeF_5 anion was reported.¹ This anion is highly unusual as it is the only known example of a pentagonal planar AX_5 species. In our search for additional representatives of this class, it was noted that the Raman spectrum of a sample² having the analytical composition Cs_3IF_6 closely resembled that of $CsXeF_5$. As the unknown IF_5^2 is isoelectronic with XeF_5 and is likely to be isostructural, we suspected that the previously reported " Cs_3IF_6 " composition might actually be a mixture of CsF and Cs_2IF_5 . Fortunately, the original sample of " Cs_3IF_6 " had been preserved in our laboratory for 25 years and showed no signs of deterioration. Therefore, this sample was reinvestigated as it could possibly contain a pentagonal planar IF_5^2 dianion. Unfortunately, both CsF and the cesium salt of the multiply

charged anion present in the " Cs_3IF_6 " sample were found to be insoluble in all available chemically inert solvents, thus preempting either their separation by extraction methods or the identification of the multiply charged anion by methods such as multinuclear NMR spectroscopy or growing of a single crystal for x-ray diffraction. It was therefore interesting to synthesize soluble IF_4 and F salts containing a common cation, which would allow the determination of the true combining ratio of IF_4 with F. A comparison of the vibrational spectra of the resulting product with those of " Cs_3IF_6 " should then also permit a positive indentification of the anion present in " Cs_3IF_6 ".

Experimental

The original ²Cs₃IF₆ sample, prepared by the combination of CsF and IF₃ in a 3:1 mol ratio in a CFCl₃ suspension at -78°C, ^{3,4} was used in this study. The preparation of IF₃, ³ anhydrous N(CH₃)₄F⁵ and XeF₂⁶ have previously been described. The purity of the IF₃ was checked by its low temperature Raman spectum which was in excellent agreement with a previous report. ⁷ The N(CH₃)₄I (K & K Labortories, Inc.) was used as received. The CH₃CN (J.T. Baker, bioanalyzed, low water) was dried over P₂O₅ prior to its use. The volatile fluorine compounds were handled in a stainless-steel Teflon-FEP vacuum line, similar to one previously described. ⁸ The CH₃CN was handled in a flamed out Pyrex glass vacuum line equipped with Kontes Teflon valves and a Heise pressure gauge.

Infrared spectra were recorded on a Perkin-Elmer Model 283 spectrophotometer using AgCl disks prepared by pressing the finely powdered sample between two thin AgCl plates in a Barnes Engineering minipress inside the glove box. An AgCl blank was placed into the reference beam to compensate for the absorption of the window material. Raman spectra were recorded on either a Cary Model 83GT, a Spex Model 1403 or a Bruker Equinox 55 spectrophotometer using the 488 nm exciting line of an Ar ion laser, the 647.1 nm line of a Kr ion laser or the 1064 nm line of a neodymium yag laser, respectively. ¹⁹F NMR spectra were recorded at 84.24 MHz on a Jeol FX902 multinuclear instrument between +30 and -30 °C, using CH₃CN as a solvent and d6-

acetone/CFCl₃ as an external lock substance and standard. The DSC data were recorded using a DuPont Model 910 DSC. A DuPont Model 2000 Thermal Analyst was used for recording and analyzing the data. The samples were crimp sealed in aluminum pans inside the drybox and heated at a rate of 2 °C/min.

Synthesis of $N(CH_3)_4IF_4$. Inside the dry box, $N(CH_3)_4I$ (2.163 mmol) and $XeF_2(4.327 \text{ mmol})$ were combined in a prepassivated (with CIF_3) Teflon FEP, 0.75 inch o.d. Teflon FEP Utube which was closed with two stainless steel valves. On the glass vacuum line, CH_3CN (10 mL liquid) was added at -196 °C, and the mixture was warmed to -31 °C. After 8 hr, all gas evolution had ceased, and the amount of gas volatile at -78 °C (2.1 mmol of Xe) was measured. A white solid had formed in the bottom of the tube which was identified as $N(CH_3)_4IF_2$. Upon warming of the mixture to room temperature, the yellowish CH_3CN solution turned more orange and additional gas evolved. After 1.5 hr at 20 °C, the gas evolution (2.13 mmol of Xe) was complete, resulting in a white solid and clear orange-yellow solution. All volatile material was pumped off at room temperature leaving behind 602 mg of a white solid (weight calcd for 2.163 mmol of $N(CH_3)_4IF_4$ = 599.4 mg) which was identified by vibrational and NMR spectroscopy as $N(CH_3)_4IF_4$.

 $N(CH_3)_4IF_4$ was also prepared by the reaction of stoichiometric amounts of $N(CH_3)_4F$ and IF_3 in CH_3CN solution of -31 °C. However, this synthesis is inferior to the one described above.

Synthesis of [N(CH₃)₄]₂IF₅. Inside the dry box, N(CH₃)₄IF₄ (0.58 mmol) and N(CH₃)₄F (0.72 mmol) were loaded into a prepassivated, 0.75 inch o.d., Teflon-FEP ampoule which contained a Teflon coated magnetic stirring bar and was closed by a stainless steel valve. On the glass line, dry CH₃CN (7.2 mL) was added at -196 °C and the mixture was stirred for 18 hr at -31 °C. The solvent was pumped off for 14 hr, first at -31 °C and during the last stages at room temperature, resulting in a white dry solid (227 mg, weight calcd for 0.58 mmol)

 $[N(CH_3)_4]_2IF_5$ plus 0.14 mmol $N(CH_3)_4F = 228$ mg) which was shown by vibrational spectroscopy to consist mainly of $[N(CH_3)_4]_2IF_5$, containing a small amount of unreacted $N(CH_3)_4IF_4$ and the excess of $N(CH_3)_4F$ used in the reaction.

Theoretical Calculations. Electronic structure calculations were done at the local density functional theory (LDFT) level^{10,11} with a polarized valence double zeta basis set (DZVP)¹², and at the Hartree-Fock (HF) level¹³ with an effective core potential (ECP)¹⁴ on I for the core electrons and with a polarized double zeta basis set for the valence electrons, and with an all electron polarized double zeta valence set on F.¹⁵ Geometries were optimized using analytical methods.¹⁶ Analytic second derivatives were calculated at the optimized LDFT geometries.¹⁷ Numerical second derivatives were calculated at the Hartree-Fock level by using a two-point differencing scheme. The LDF calculations were done with DGauss¹⁸, and the HF calculations were done with Gaussian 94.¹⁹ The calculated Hessian matrices (second derivatives of the energy with respect to Cartesian coordinates) were converted to symmetry-adapted internal coordinates for further analysis using the program systems GAMESS²⁰ and Bmtrx.²¹

Results and Discussion

Synthesis and Properties of $N(CH_3)_4IF_4$. As already pointed out in the Introduction, the determination of the combining ratio of IF_4 with F required the availability of soluble salts of these two anions with a common counterion. Since $N(CH_3)_4F^5$ and $N(CH_3)_4XF_4$ (x = Br or Cl)²² were known to have good solubility in CH_3CN , the $N(CH_3)_4IF_4$ salt appeared to be an ideal candidate for our reactions. Our first attempts to prepare this salt from $N(CH_3)_4F$ and IF_3 in CH_3CN solution at -31 °C produced only impure products. Subsequently, the following improved synthesis (1) gave very pure $N(CH_3)_4IF_4$, eliminated the need for the thermally unstable IF_3 , which is difficult to purify,^{3,4} and utilized only commercially available starting materials. It is

$$N(CH_3)_4I + 2XeF_2$$
 CH_3CN $CH_3)_4IF_4 + 2Xe$ (1)

similar to the one previously reported 23 for 1,1,3,3,5,5 hexamethylpiperidinium (pip⁺)IF₄ from pip⁺F⁻, iodine and XeF₂, but offers the advantages of avoiding a 7 step, 14% yield syntesis of pip⁺F and the handling of volatile elemental iodine.

N(CH₃)₄IF₄ is a white crystallinic solid which, based on DSC data and vibrational spectra, undergoes a reversible endothermic phase change of 21.8 J/g at 73.8 °C and irreversible exothermic decomposition at 284 °C. The surprisingly high thermal stability of this compound is in marked contrast to that of pip⁺IF₄ which was reported²³ to be unstable above 0 °C.

The ¹⁹F NMR spectra of N(CH₃)₄IF₄ in CH₃CN solution were recorded at -30 and 30 °C and consisted of broad singlets at δ = -106.5 with half widths of 385 and 1061 Hz, respectively. The observed chemical shift is in excellent agreement with the value of -106 ppm, previously reported²³ for pip⁺IF₄ in CH₃CN at -20 °C. The chemical shift of IF₄ fits nicely the trends observed for other halogen fluoride anions, ^{25,26} i.e., the shielding of the fluorine ligands decreases with increasing oxidation state and electronegativity of the central atom (see Table 1). Pronounced line broadening with increasing temperature was also observed for BrF₄ and ClF₄. ²⁶

The infrared and Raman spectra of $N(CH_3)_4IF_4$ are summarized in Table 2. The bands observed for the IF_4 part of $N(CH_3)_4IF_4$ agree well with those previously reported for $CsIF_4$, but not with those (Ra: 573.5, 461 and 213 cm⁻¹) previously listed²³ for pip^+IF_4 . Particularly, the value of 573.5 cm⁻¹, reported for $v_1(A_{1g})$ of pip^+IF_4 , is outside the expected (see Computional Results below) range and is likely incorrect. The bands due to the $N(CH_3)_4$ cation have been summarized in a footnote of Table 2. They are in excellent agreement with those previously reported for other $N(CH_3)_4$ salts^{1.5,9} and, therefore, do not require further discussion.

Synthesis and Properties of $[N(CH_3)_4]_2IF_5$. The combination of $N(CH_3)_4IF_4$ with a slight excess of $N(CH_3)_4F$ in CH_3CN solution at -31 °C produced the desired $[N(CH_3)_4]_2IF_5$ (2).

$$N(CH_3)_4IF_4 + N(CH_3)_4F \xrightarrow{CH_3CN} [N(CH_3)_4]_2IF_5$$
 (2)

Prolonged reaction times (about 20 hr) and stirring resulted in an almost quantitative conversion of IF_4 to IF_5^{2} thus establishing the 1:1 combining ratio between IF_4 and F. In the presence of a large excess of F ions, no further fluoride ion uptake and, hence, no formation of IF_6^{3} were observed. Attempts to carry out reaction (2) in CHF₃ solution at -78 °C were unsuccessful and only unreacted $N(CH_3)_4IF_4$ was recovered.

The $[N(CH_3)_4]_2IF_5$ salt is a white solid which is stable at room temperature. It has very little solubility in solvents, such as CH_3CN , which precluded its characterization by NMR spectroscopy and the growing of single crystals for a crystal structure determination. Infrared and Raman spectra of the solid were used for its characterization. They are given in Figure 1 and Table 3, and their assignments are discussed below in more detail.

Characterization of Cs_2IF_5 . In the original report² on " Cs_3IF_6 ", the true combining ratio of IF_3 with CsF had not been established. An arbitrary 1:3 ratio of the starting materials had been used and, since both, CsF and Cs_2IF_5 , are insoluble, the products could not be separated and, therefore, by necessity had the analytical composition Cs_3IF_6 .

In the reaction of IF_4^{-1} with an excess of fluoride ions, the logical first step is the addition of one fluoride ion with formation of IF_5^{-2} . The addition of a second fluoride ion with formation of an IF_6^{-3} trianion becomes less likely since the Lewis acidity of the parent anion decreases with the addition of each F^- ion. Although, based on these arguments, the formation of IF_5^{-2} is clearly favored over that of IF_6^{-3} , it previously had not been given much consideration due to the absence of any other AX_5E_2 species containing two free valence electron pairs (E) on A. This picture, however, has changed recently with the synthesis of surprisingly stable XeF_5^- salts.¹

Our reinvestigation of the original² "Cs₃IF₆" sample, which had been preserved in our laboratory for 25 years without noticeable decomposition, showed that is it is indeed a mixture of Cs_2IF_5 and CsF, based on the following evidence:

- (i) X-ray powder data: The previously published powder pattern of "Cs₃IF₆" (Table II of reference 2) shows all the lines characteristic²⁷ for CsF (d(Å), intens: 3.4 ms, 3.03 mw, 2.106 mw, 1.801 mw, 1.735 vw, 1.496 w, 1.370w, 1.336 mw, 1.223 w, 1.153w) and leaves no doubt that "Cs₃IF₆" contains a very significant amount of free CsF.
- (ii) Infrared spectrum: The infrared spectrum of "Cs₃IF₆" was rerecorded and compared to that of CsF taken under the same conditions. It was found that the major absorption in both spectra was a very intense and broad band at 298 cm⁻¹ due to CsF which in the "Cs₃IF₆" case confirms the presence of free CsF but obscures the less intense IF₅² bands.
- (iii) Raman spectrum: A higher quality Raman spectrum, which did not suffer from fluorescence,² was obtained by means of 1064 nm excitation with a neodymium yag laser and is shown in Figure 2. The dominant Raman bands at 478, 396, 339 and 325 cm⁻¹ are in excellent agreement with those of IF_5^2 in $[N(\text{CH}_3)_4]_2 \text{IF}_5$ (see Table 3) and leave no doubt that the dominant anion in " Cs_3IF_6 " is IF_5^2 . For a hypothetical IF_6^3 anion, the additional fluoride ligand and extra formal negative charge should cause a different band pattern and significant shifts to lower frequencies. In addition to the intense IF₅² bands, the "Cs₃IF₆" spectrum showed several very weak bands which can be assigned to likely impurities, such as IF_7^{2-} (556 and 454 cm⁻¹) or IOF_5^{2-} (865 cm⁻¹)²⁸ and lattice vibrations (140 and 110 cm⁻¹). The formation of IF₇² and IOF₅² can be easily rationalized. Iodine trifluoride has a tendency to disproportionate to I₂ and IF₅, and IF₅ is known²⁸ to react with excess F to give IF₇². Similarly, a trace of moisture will result in partial hydrolysis of IF₆ to give IOF₄, which in turn is known to give with excess F the IOF₅²⁻ dianion.²⁸ These impurities had been formed already during the original preparation⁴ and were not generated during sample storage.

Vibrational Spectra and Electronic Structure Calculations. The vibrational spectra of IF_5^{2-} are summarized in Table 3 and were assigned by analogy with those established for the isoelectronic XeF_5^{-} anion.^{1,29} The only difference is the reversal of the identity of ν_5 amd ν_6 in the E_2^{-} block, based on the results of the potential energy distribution (see below). The spectra of IF_5^{-2-} and

 XeF_5 are very similar. They exhibit the expected frequency decreases for IF_5^{2-} due to the additional negative charge which enhances the $X(\delta+)$ - $F(\delta-)$ polarity of the bonds. Analogous general frequency and bond weaking effects are also observed on going from IF_4 (see Table 2) to IF_5^{2-} (see Table 3).

The geometry and the vibrational spectra of IF_5^{2} and XeF_5 were studied by electronic structure calculations at the HF and LDFT levels. To evaluate the reliability of our computational methods, the well know IF_4 geometry²³ and spectra² were calculated first. As can be seen from Tables 2 and 4, the HF/ECP/DZP method gave the better results. For IF_4 , the bond length, calculated at the HF/EDP/DZP level is only 0.013 Å shorter than the observed one.²³ For IF_5^{2-} , the minimum energy structure at both the HF and LDFT levels was pentagonal planar with D_{5h} symmetry. Assuming a similar correction as for IF_4 , the bond length of IF_5^{2-} is predicted to be 2.11 Å. This bond length increase of 0.09 Å relative to IF_4 can be explained by the increased polarity of the bonds (see above) and increased ligand-ligand repulsion due to a decrease of the F-I-F bond angle from 90 ° in IF_4 to 72 ° in IF_5^{2-} .

The symmetry force constants and potential energy distributions (PED) were calculated for $\text{IF}_5^{\ 2}$ and XeF_5 using the scaled HF/ECP/DZP frequencies (see Table 5). Since we were unable to duplicate with our computer calculations the previously published, hand calculated, out of plane deformation force constants, $^{1.29}$ the originally given 1 G matrix was reexamined and found to contain incorrect multiplicity factors for G_{22} and G_{77} . The corrected G matrix for XeF_5 is given in Table 6 and was verified by the machine methods. A typographical error for a sign in one of the originally published symmetry coordinates, i.e., (S_{5b}) should read $(2/5)^{1/2}$ [$\sin 2\alpha(\Delta r_2 - \Delta r_5) - \sin\alpha(\Delta r_3 - \Delta r_4)$] has already been corrected elsewhere. The revised force constant, $F_{22} = 0.480$ mdyn Å/rad², for the symmetric out of plane deformation of XeF_5 is now in much better agreement with our expectation. The deformations out of the highly crowded pentagonal plane should possess significantly smaller force constants than the corresponding in plane deformations.

As can be seen from Tables 3 and 5, the results from the normal coordinate analyses confirm the identities of these pentagonal planar anions. The general fit between the calculated and observed frequencies is very good. The only remaining minor ambiguities in the normal coordinate analyses are: (i) for XeF_5 , the exact location of the infrared active, in plane deformation mode, v_4 , is somewhat uncertain. Based on the calculations, its frequency is similar to, but its infrared intensity is much lower than those of the out of plane umbrella deformation mode, v_2 . Therefore, we assume that v_4 is hidden under the v_2 band; (ii) for $\mathrm{IF_5}^{2}$, the location of the infrared active, out of plane umbrella deformation mode was obscured. In the Cs₂IF₅ spectrum, the strong absorption due to free CsF obscured this region, and in the $[N(CH_3)_4]_2IF_5$ spectrum, the presence of an IF_4 impurity and the broadness of the v_3 mode of ${\rm IF_5}^2$ interferred. In the Raman spectrum of ${\rm Cs_2IF_5}$, one of the two E2 bands exhibits a splitting, similar to those observed for these modes in XeF5. In Table 3, this splitting was tentatively assigned to v_5 because this requires a smaller splitting. However, if one assumes v_6 to be split intead of v_5 , the average of the 396 and 339 cm⁻¹ components would result in much better agreement with the calculated frequency of 366 cm⁻¹ and the [N(CH₃)₄]₂IF₅ spectrum, in which v_6 of IF₅²⁻ probably coincides with a cation band at 366 cm⁻¹.

An inspection of the potential energy distributions for ${\rm IF}_5^{2^-}$ and ${\rm XeF}_5^-$ (see Table 5) shows that the E modes of ${\rm XeF}_5^-$ are considerably more characteristic than those of ${\rm IF}_5^{2^-}$. This is due to the increased bond polarity in ${\rm IF}_5^{2^-}$, which lowers preferably the stretching force constant and makes its value more similar to those of the deformation constants. This enhances the mixing of the normal modes in ${\rm IF}_5^{2^-}$ and, in its ${\rm E}_2^+$ block, results in the higher frequency mode becoming mainly the bending motion. A further consequence of the larger bond polarity in ${\rm IF}_5^{2^-}$ is the increased coupling of the stretching motions involving opposite bonds (frr' = 0.42 mdyn / Å). This large frr' value is responsible for the low frequency of the antisymmetric stretching mode , ${\rm v}_3({\rm E}_1^+)$, which respresents

the reaction coordinate for the loss of an F ion. This finding is in accord with the observation that IF_5^2 readily loses an F ion to give IF_4 .

The stability of the hypothetical ${\rm IF_6}^{3-}$ anion was also examined by ab initio methods. The four most likely geometries of ${\rm IF_6}^{3-}$ are shown in Figure 3. Like ${\rm IF_5}^{2-}$, this anion contains two free valence electron pairs on iodine. If both electron pairs are sterically active, the most likely structure is the one with ${\rm D_{3d}}$ symmetry, in which the two more repulsive free pairs avoid each other as much as possible by occupying the axial trans positions of a bicapped octahedron. If, on the other hand, one of the two free valence electron pairs on iodine is sterically inactive, i.e., occupies an ${\rm A_{1g}}$ or sorbital, only p and d-orbitals of iodine are involved in the bonding and, therefore, the resulting structures are governed by repulsion effects, and, as in transition metal heptafluorides, the energetically most favored structures become the monocapped octahedron of ${\rm C_{3v}}$ symmetry and the monocapped trigonal prism of ${\rm C_{2v}}$ symmetry. The pentagonal bipyramid of ${\rm C_{5v}}$ symmetry is slightly higher in energy, but was also considered.

The structures of all four geometries of ${\rm IF_6}^{3-}$ were calculated at the restricted Hartree-Fock (RHF) self-consistent field level, using effective core potentials³⁴ and the corresponding valence basis sets of Stevens, Basch, Krauss, and Jasien.³⁵ The basis set was augmented with a diffuse s+p shell³⁶ and a d polarization function³⁷ on each atom. All calculations were performed using the GAMESS quantum chemistry package.²⁰ Since the HOMO of the D_{3d} and C_{3v} structures is a half-filled degenerate orbital pair, Jahn-Teller distortion to lower symmetry is anticipated. Therefore, the initial geometries of these conformations were distorted to C_s symmetry by slight elongation of one pair of trans fluorine ligands. For each of the four conformations, geometry optimization led to the dissociative loss of one or two fluoride ligands. Specifically, optimization of the C_{2v} conformation and the distorted D_{3d} and C_{3v} structures led to formation of square planar $[{\rm IF}_4]^2$ and 2 fluoride anions, whereas optimization of the C_{5v} structure led to formation of pentagonal planar $[{\rm IF}_5]^{2v}$ and one fluoride anion. Thus, none of the four conformations were found to be a local minimum. Although more extensive calculations are desirable to definitively rule out the stability of the $[{\rm IF}_6]^3$, the present results strongly indicate that this anion is not a stable species.

Conclusions

The successful synthesis of a second example of a pentagonal planar AX_5 species demonstrates that XeF_5 is not a unique case, and that other examples of AX_5 molecules of D_{5h} symmetry, such as TeF_5^{3-} , AuF_5^{2-} or PtF_5^{3-} , might also exist. Although IF_5^{2-} had been prepared 30 years ago,³ its nature had not been recognized until now. The greatly delayed recognition of this unusual anion parallels the discovery¹ of XeF_5^{-} which had originally been mistaken for octahedral XeF_6^{2-} .³⁸⁻⁴⁰

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- Dedicated to Professor Neil Bartlett on the occasion of his 65th birthday
- ‡ Hughes STX
- § University of Southern California
- || Pacific Northwest National Laboratory
- # Propulsion Directorate, Air Force Research Laboratory
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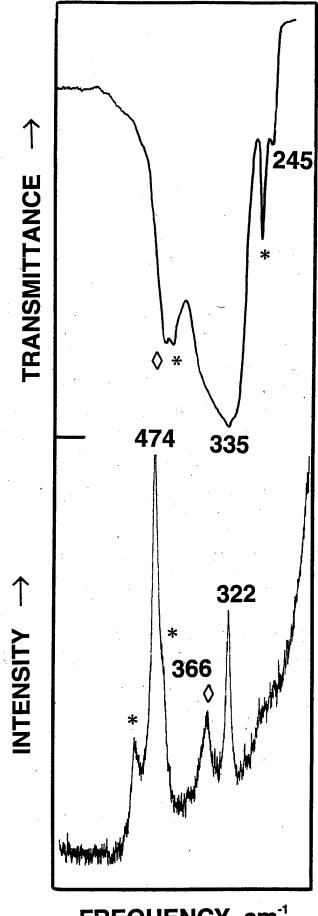
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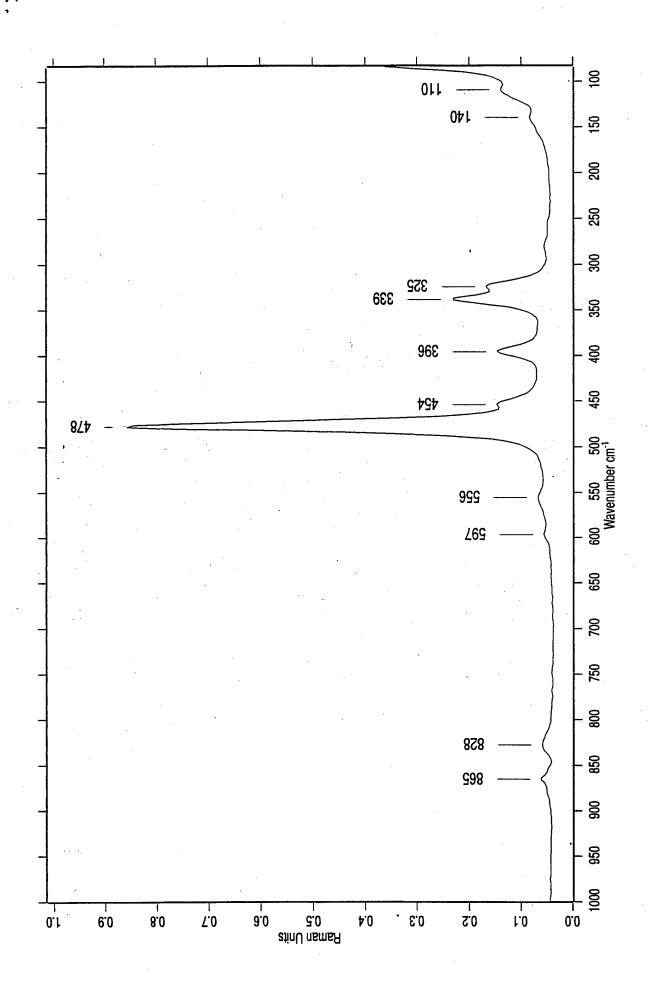
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Diagram Captions

Figure 1	Infrared and Raman spectra of solid $[N(CH_3)_4]_2$ IF ₅ . The bands masked by
.·	asterisks and diamonds are due to IF_4 and $N(CH_3)_4^+$, respectively, while those
	marked by frequency values belong to IF ₅ ² .
Figure 2	Raman spectrum of Cs ₂ IF ₅ .
Figure 3	Most likely geometries of the IF ₆ ³ trianion and their spontaneous decomposition
	modes.

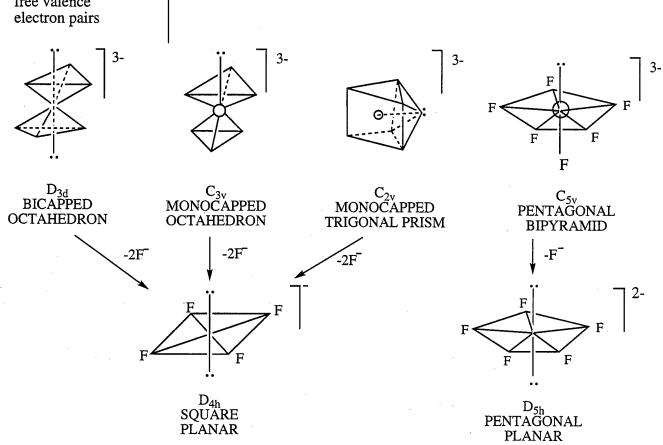


FREQUENCY, cm⁻¹



Two sterically active free valence electron pairs

One sterically active and one sterically inactive free valence electron pair



BrF₆.

Table 1. 19F Chemical Shifts for Binary Halogen Fluoride Anions

oxidation state of central atom +I +III +VII	IF282* IF106* IF -
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^a. ref 9. ^b this work and ref 23. ^c ref 25. ^d ref 26. ^c ref 22.

Table 2. Observed and Calculated Vibrational Spectra of IF

assgnts descrip	and a	pproy n poi	assgnts and approx mode descriptions in point group D _{4h}	CsIF	obsd freq, cm ⁻¹ (rel int) ^a $CsIF_4 \qquad N(CH_3)_4IF_4^4$ $RA \qquad IR$	m ⁻¹ (rel int) N(CH ₃) IR	, IF d RA	calcd fre HF/ECP/DZP unscaled scaled	calcd freq, cm ⁻¹ (IR int) ^b P/DZP LD] I scaled uns	nt) ^b LDFT/DZVP unscaled
Ra	$\mathbf{A}_{i\mathbf{g}}$		v ₁ v sym in phase	. 1	522(10)	1	515(10)	559(0)	520	495(0)
띪	A_{2u}		v ₂ ô umbrella	271ms	1,	267ms	·	326(78)	772	255(40)
Ra	$\mathbf{B}_{1\mathbf{g}}$		v ₃ v sym out of phase	1	455(7.2)	I	457(5.6)	493(0)	458	449(0)
Ra	$\mathbf{B}_{2\mathbf{g}}$		V ₄ δ scissor	1	195(0+)	I	197(0.5)	225(0)	191	167(0)
l	$\mathbf{B}_{2\mathbf{u}}$	>,	v _s ð pucker	1		·		185(0)	157	140(0)
IR	ញ	>9	v ₆ v asym	448vs	1	449 vs		481(951)	447	501(544)
		Y	v, 8 asym	n. obsd	1 .	n. obsd		139(3)	118	115(3)

^aData for CsIF₄ from ref 2, for N(CH₃)₄ IF₄ from this study. b Infrared intensities in km/mol. Empirical scaling factors of 0.9294 and 0.8490 addition to the above listed IF₄ bands, the following bands due to N(CH₃)₄ were observed: IR: 3118sh, 3040w, 2969vw, 1490 mw, 1444w, 1416mw, 1287w, 952m, 922w, 462m; Ra: 3021(0.7), 2969(0.2), 2942(0.4), 2907(0.1), 2800(0.1), 1471(0.6), 1461(0.7), 1413(0.3), for the stretching and deformation modes, respectively, were used to maximize the fit between the observed and calculated frequencies. In 1284(0.1), 1175(0.15), 1168(0.1), 947(1.5), 754(1.5), 450 sh on the intense 457 IF₄ band, 369(0.2)

Table 3. Observed and Calculated Vibrational Spectra of IF52 Compared to Those Observed for XeF5

	×	Ra	502(10)				423(2.1)		377(3.3)	
	r (rel int N(CH	R		278s	509sh 465vs 420sh					
XeF.*	obsd fr	IK Ka	504(10)	274s	450vs 415s	288sh	432(1.5)	422(1.6)	380(2.2) 369(2.3)	
	m ⁻¹ (IR int) ^f //DZP	scaled	468	307	342	249	335		366	100
	calcd freq, cm ⁻¹ (IR int) ^f HF/ECP/DZP	unscaled	487(0)	320(97)	356(588)	259(14)	349(0)		381(0)	104(0)
IF 2.				•	÷ ,					
	obsd freq, cm ⁻¹ (rel int) Cs ₂ IF ₅ [N(CH ₃) ₄] ₂ IF ₅ ^b RA IR DA	WW WI	474(10)	(—) _e	335vs,br	245w	322(3.8)) (366)	
	obsdí Cs ₂ IF ₅ RA	54	478(10)				339(2)	325(1.6)	396(0.9)	
	assgnts and approx mode descriptions in point group D _{ss}		v sym	ð umbrella	v asym	8 asym in plane	v _s v asym		ð asym in plane	v, 8 pucker
	pprox r in point		<u>۲</u>	^	× ×	>	^		>	٧,
	assgnts and approx mode descriptions in point grou		A ₁ ′	A ₂ ′′	ъ		'n			n,
	assgridescr		Ra	R	Ħ		Ra			Ra

*Data from ref 1. b In addition to the above listed IF₅²⁻ bands, the following bands due to N(CH₃)₄⁺ were observed: IR: 3034ms, 1507ms, 1415w, 1255m, 963s, 467m; Ra: 3010(1.3), 2943(0.8), 2810(0.4), 1478(1.8), 952(1.8), 751(3.0), 460(sh), 366(1.0). ° This band is masked by the very intense broad band at 335 cm⁻¹. This band is obscured by the relatively intense 366 cm⁻¹ N(CH₃)₄⁺ band. °An empirical scaling factor of 0.9606 was used to maximize the fit between observed and calculated frequencies. Ir intensities in km/mol.

Table 4. Calculated, Observed and Predicted Bond Lenghts for IF, and IF, 2-

(Å)nted predicted	LDFT/DZVP	2.069	2.163 2.11
——— bond distances (Å) calculated	HF/ECP/DZP	1.994	2.094
observed		2.007	1
		$F_4(D_{4h})$	$F_{5}^{2-}(D_{5h})$

^a Ref. 23.

Table 5. Symmetry Force Constants and Potential Energy Distribution of D_{5h} IF₅² and XeF₅ Calculated from the Scaled HF/ECP/DZP Second Derivatives

			IF ₅ ^{2-a}		-		XeF ₅ -b	
freq, cm ⁻¹	m.i		sym force consts°	PED (%)	freq, cm ⁻¹	m-1	sym force consts ^c	PED (%)
	psqo	calcd			psqo	calcd		
\mathbf{A}_{l}	474	468	$F_{11} = 2.445$	100(1)	502	503	$F_{11} = 2.835$	100(1)
Α,"		307	$F_{22} = 0.529$	100(2)	278	305	F., = 0.480	100(2)
, П	335	342	$F_{33} = 0.867$	70(3),20(4)	455	454	$F_{33} = 1.707$	98(3),2(4)
ŗ	245 322/339	249 335	$F_{34} = -0.030$ $F_{44} = 1.886$ $F_{55} = 1.735$	80(4),20(3) 74 (5),26(6)	278 423	276 425	$F_{34} = -0.181$ $F_{44} = 1.952$ $F_{55} = 1.982$	98(4),2(3) 92(5),8(6)
, <u>ந</u>	366/396	366	$F_{56} = 0.115$ $F_{66} = 1.316$ $F_{77} = 0.260$	74(6),26(5) 100(7)	377	378	$F_{56} = 0.122$ $F_{66} = 1.794$ $F_{77} = 0.266$	92(6),8(5) 100(7)
internal	internal force constants	nts			·	•		
		£ ££	1.529 0.035 0.423				2.043 0.137 0.260	

*Empirical scaling factors of 0.9606 and (0.9606)² were used for the frequencies and force constants, respectively. ^bScaling factors: 0.87322 and (0.87322)². Stretching constants in mdyn/Å, deformation constants in mdyn Å/rad², and stretch-bend interaction constants in mdyn/rad.

Table 6. G Matrix* for Pentagonal Planar XeF, of Symmetry D,

$$\lambda_1' \qquad G_{11} = \mu_y = 5.2637 \times 10^{-2}$$

$$\lambda_{2}'' \qquad G_{22} = (5/r^{2})(\mu_{y} + 5\mu_{x}) = 1.1201 \times 10^{-1}$$

$$G_{33} = \mu_y + 5u_x/2 = 7.1679 \times 10^{-2}$$

$$G_{34} = (5^{3/2}\mu_x)/(4r\sin\alpha) = 1.1123 \times 10^{-2}$$

$$G_{44} = (1/r^2)(5\mu_y \sin^2 2\alpha + \mu_x) = 2.4334 \times 10^{-2}$$

$$G_{55} = \mu_y = 5.2637 \text{ x } 10^2$$

Ę,

$$G_{56} = 0$$

$$G_{66} = (1/r^2)(4\mu_y \sin^2 \alpha) = 4.7026 \text{ x } 10^{-2}$$

$$E_2'' G_{77} = (3-5^{1/2})(5/2r^2)\mu_y = 2.4823 \times 10^{-2}$$

^{*}The following geometry and masses were used for the calculation of the G matrix: r = 2.0124 Å and $\alpha = 72^{\circ}$; $m_x = 131.292$; $m_y = 18.998$.